

SST $k - \omega$ Simulations of the Atmospheric Boundary Layer Including the Coriolis Effect

Christiane Adcock¹, Marc Henry de Frahan², Jeremy Melvin³, Ganesh Vijayakumar², Shreyas Ananthan⁴, Gianluca Iaccarino¹, Robert Moser³, Michael Sprague²

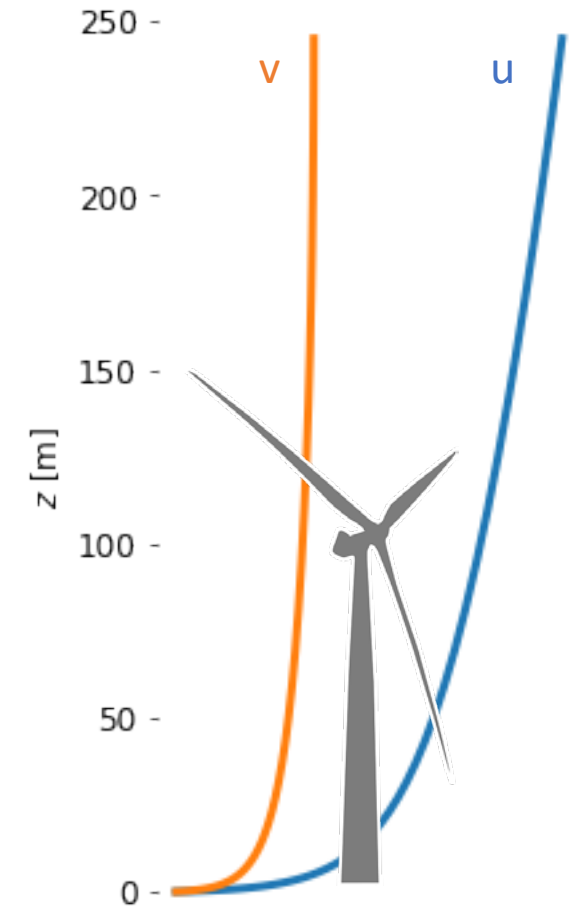
¹Stanford University ²National Renewable Energy Laboratory ³University of Texas at Austin ⁴Siemens Gamesa Renewable Energy

Blade-resolved simulations of wind turbines in the Atmospheric Boundary Layer (ABL) [1]

Most ABL simulations use LES or RANS $k - \varepsilon$

- LES prohibitively expensive for blade-resolved + multiple turbines [1]
- RANS $k - \varepsilon$ inaccurate for adverse pressure gradients near blades [2,3]
- RANS $k - \omega$ overly sensitive to freestream value of ω [2,3]

Current goal: RANS SST $k - \omega$ to simulate ABL w/ Coriolis, w/out buoyancy, w/out turbines

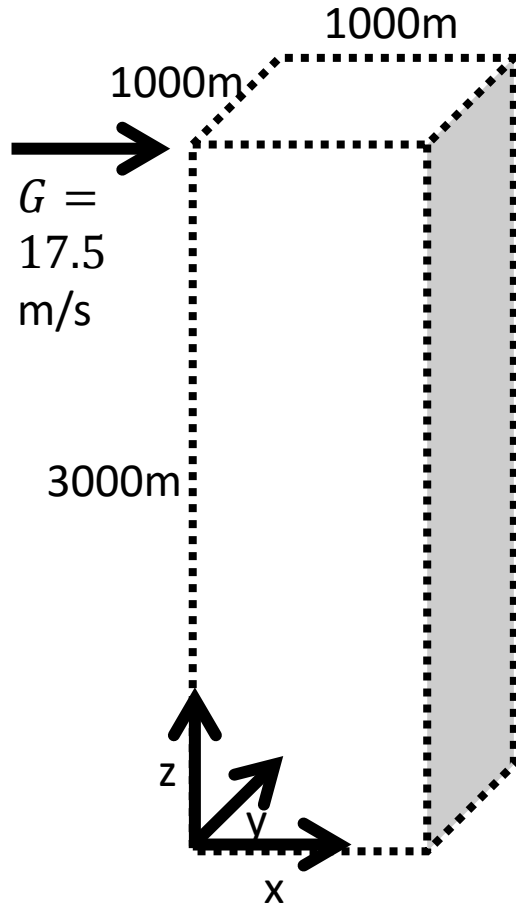


[1] 2020 Sprague et. al, NAWEA WindTech

[2] 1992 Menter, NASA Technical Memorandum

[3] 2006 Wilcox, Turbulence modeling for CFD

Test Case Set-Up



BCs: Top=symmetry. Sides=periodic. Bottom:

- $u = 0. k = \frac{u_\tau^2}{\sqrt{\beta^*}}; \omega = \frac{u_\tau}{\sqrt{\beta^* \kappa z_0}}$ [4]
- Derive: $\varepsilon = \frac{u_\tau k \sqrt{\beta^*}}{\kappa z_0}$ by $\omega = \frac{\varepsilon}{\beta^* k}$

Model constants: SST $k - \omega$ [4]; $k - \varepsilon$ [5]

Implemented in **Nalu-Wind**: open-source, incompressible, massively parallel flow solver for simulations of turbines in wind farms. <https://github.com/Exawind/nalu-wind>

Test cases:

- (1) W/out Coriolis. (2) W/ Coriolis:

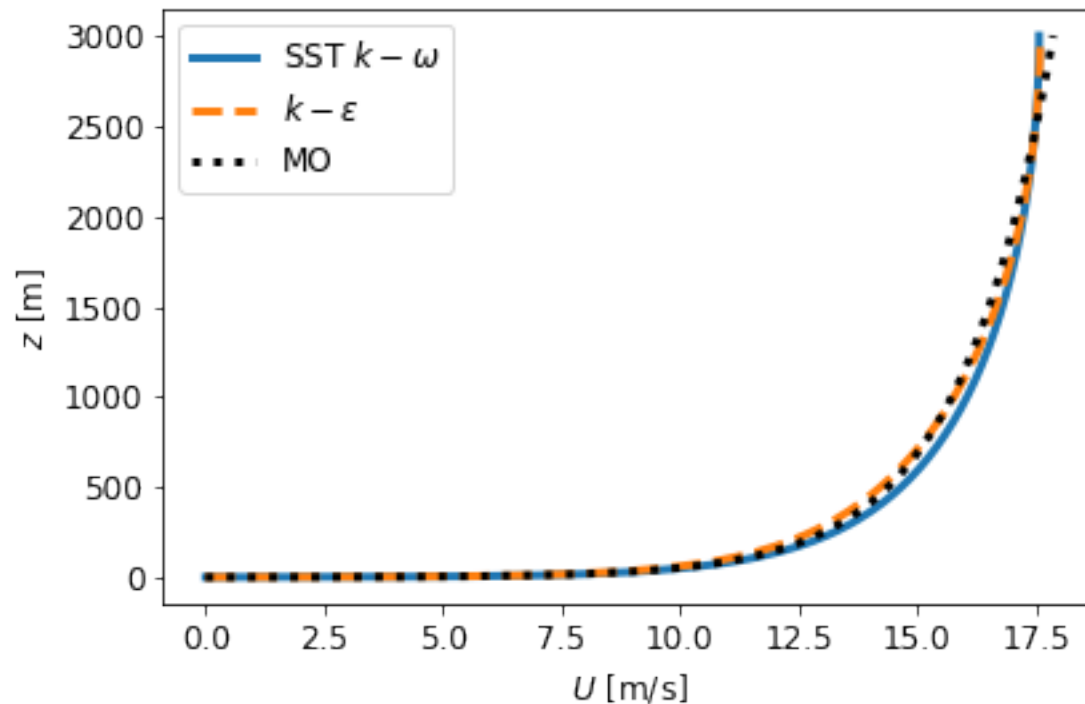
$$\frac{\partial U_i}{\partial t} + \frac{\partial U_i U_j}{\partial x_j} = -\frac{1}{\rho} \frac{\partial P}{\partial x_i} + \frac{\partial (2\nu S_{ij})}{\partial x_j} + \frac{\partial \tau_{ij}}{\partial x_j} + F_i^C$$

[1] 2020 Sprague et al, NAWEA WindTech

[5] 2013 Koblitz, Ph.D. thesis, DTU

[4] 2015 Bautista, Dufresne, Masson, E3S Web of Conferences

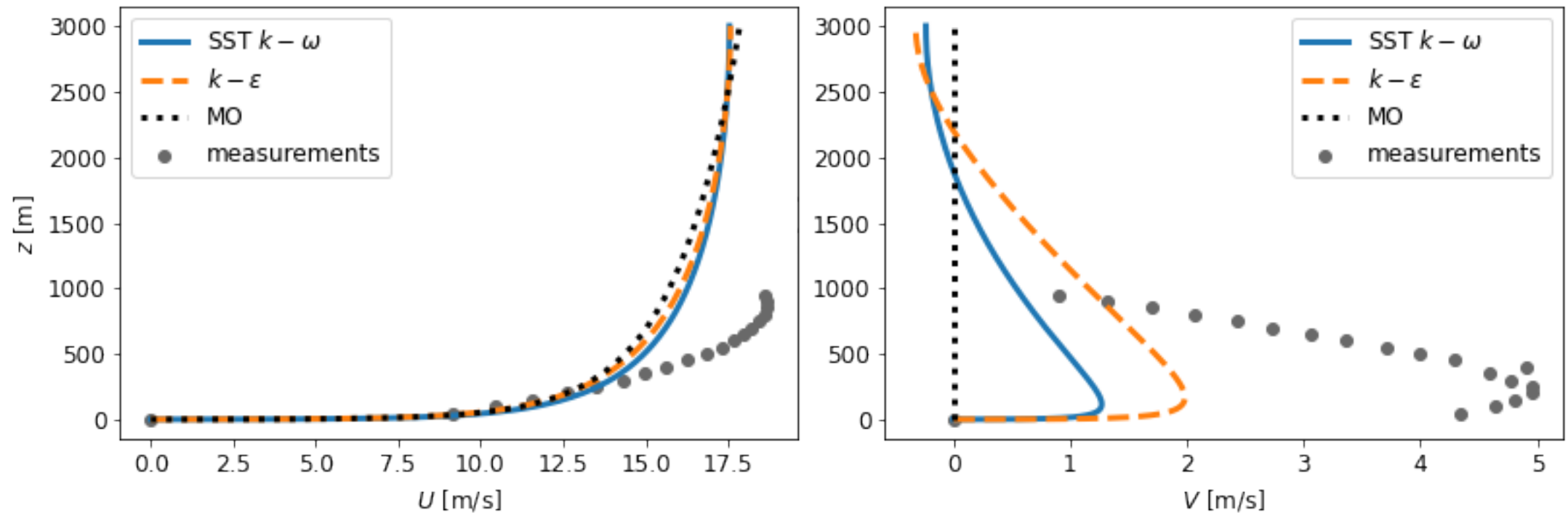
Baseline SST $k - \omega$ and $k - \epsilon$ w/out Coriolis match Monin-Obukhov similarity theory



- Monin-Obukhov (MO) log wind profile:

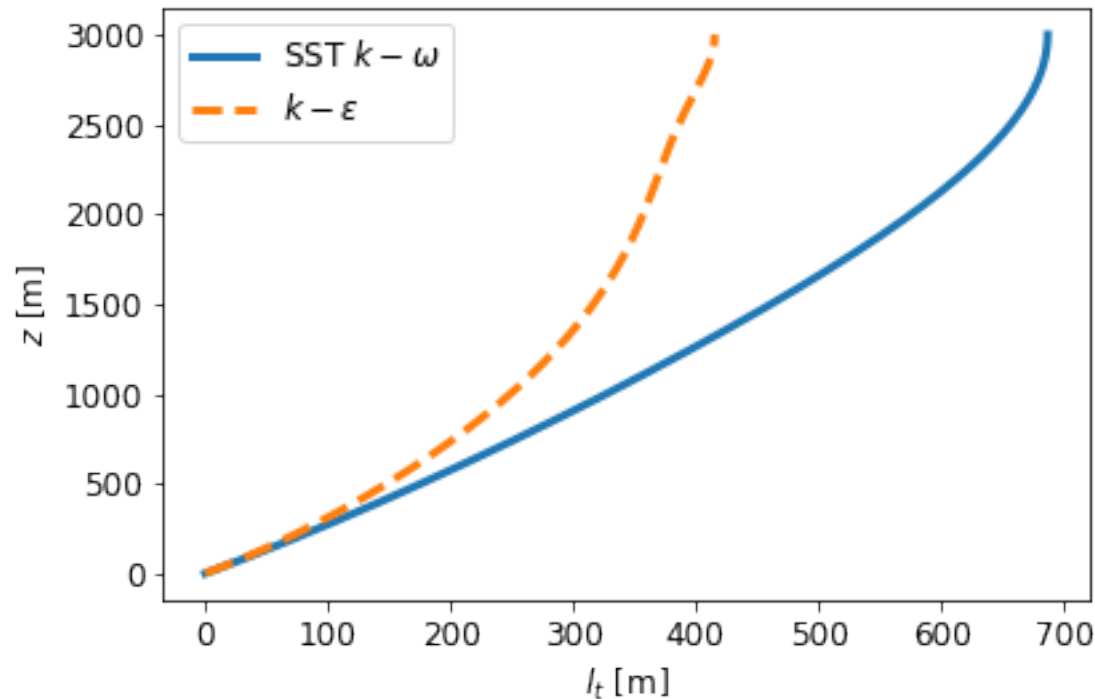
- $$u = \frac{u_\tau}{\kappa} \ln \left(\frac{z - z_0}{z_0} \right)$$

Baseline SST $k - \omega$ and $k - \epsilon$ w/ Coriolis don't fully capture Coriolis effect



Measurements from Leipzig field test [6]

Baseline SST $k - \omega$ and $k - \epsilon$ don't capture Coriolis effect b/c of large mixing length



SST $k - \omega$ blending function, F_1 :

- 0 away from wall ($k - \epsilon$)
- 1 near wall ($k - \omega$)

Here, $F_1 \approx 1$ always b/c in F_1 :

$$\tanh \left(\left(\max \left(\frac{l_t}{d}, \frac{\cdot}{d^2} \right) \right)^4 \right)$$

$$= \tanh(\text{large}) \approx 1$$

$k - \varepsilon$ mixing length limiter [5]

$k - \varepsilon$ [7]: RANS, eddy viscosity model, $\nu_t = C_\mu \frac{k^2}{\varepsilon}$, 2 eqn. — k, ε

- ε eqn.:
$$\frac{\partial \varepsilon}{\partial t} + U_j \frac{\partial \varepsilon}{\partial x_j} = \frac{\varepsilon}{k} (C_{\varepsilon 1} P - C_{\varepsilon 2} \varepsilon) + \frac{\partial}{\partial x_j} \left[\left(\nu + \frac{\nu_T}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right]$$

$k - \varepsilon$ mixing length limiter [5]

$k - \varepsilon$ [7]: RANS, eddy viscosity model, $\nu_t = C_\mu \frac{k^2}{\varepsilon}$, 2 eqn. — k, ε

- ε eqn.:
$$\frac{\partial \varepsilon}{\partial t} + U_j \frac{\partial \varepsilon}{\partial x_j} = \frac{\varepsilon}{k} (C_{\varepsilon 1} P - C_{\varepsilon 2} \varepsilon) + \frac{\partial}{\partial x_j} \left[\left(\nu + \frac{\nu_T}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right]$$

$C_{\varepsilon 1} \rightarrow C_{\varepsilon 1}^* = C_{\varepsilon 1} + (C_{\varepsilon 2} - C_{\varepsilon 1}) l_t / l_e$

- Mixing length: $l_t = \beta^{*3/4} k^{3/2} / \varepsilon$
- Maximum mixing length: $l_e = 0.00027 G / f_c$
- Coriolis force: $f_c = 2\Omega \sin \lambda$

[5] 2013 Koblitz, Ph.D. thesis, DTU

[7] 2003 Menter et. Al., Turbulence, Heat and Mass Transfer

$k - \varepsilon$ mixing length limiter [5]

$k - \varepsilon$ [7]: RANS, eddy viscosity model, $\nu_t = C_\mu \frac{k^2}{\varepsilon}$, 2 eqn. — k, ε

- ε eqn.: $\frac{\partial \varepsilon}{\partial t} + U_j \frac{\partial \varepsilon}{\partial x_j} = \frac{\varepsilon}{k} (C_{\varepsilon 1} P - C_{\varepsilon 2} \varepsilon) + \frac{\partial}{\partial x_j} \left[\left(\nu + \frac{\nu_T}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right]$

$$C_{\varepsilon 1} \rightarrow C_{\varepsilon 1}^* = C_{\varepsilon 1} + (C_{\varepsilon 2} - C_{\varepsilon 1}) l_t / l_e$$

- Mixing length: $l_t = \beta^{*3/4} k^{3/2} / \varepsilon$

- Maximum mixing length: $l_e = 0.00027 G / f_c$

- Coriolis force: $f_c = 2\Omega \sin \lambda$

As $\frac{l_t}{l_e} \rightarrow 1$, $C_{\varepsilon 1}^* \rightarrow C_{\varepsilon 2}$ so $\frac{\varepsilon}{k} (C_{\varepsilon 1}^* P - C_{\varepsilon 2} \varepsilon) \rightarrow C_{\varepsilon 2} \frac{\varepsilon}{k} (P - \varepsilon)$ so limit l_t

As $\frac{l_t}{l_e} \rightarrow 0$, $C_{\varepsilon 1}^* \rightarrow C_{\varepsilon 1}$ so $\frac{\varepsilon}{k} (C_{\varepsilon 1}^* P - C_{\varepsilon 2} \varepsilon) \rightarrow \frac{\varepsilon}{k} (C_{\varepsilon 1} P - C_{\varepsilon 2} \varepsilon)$ so reverts

[5] 2013 Koblitz, Ph.D. thesis, DTU

[7] 2003 Menter et. Al., Turbulence, Heat and Mass Transfer

Derivation of SST $k - \omega$ length scale limiter

SST $k - \omega$ [8]: RANS, eddy viscosity model, $\nu_t \approx \frac{k}{\omega}$, 2 eqn.— k, ω

- ω eqn.:
$$\frac{\partial(\rho\omega)}{\partial t} + \frac{\partial(\rho u_j \omega)}{\partial x_j} = \frac{\gamma}{\nu_t} P - \beta \rho \omega^2 + \frac{\partial}{\partial x_j} \left[(\mu + \sigma_\omega \mu_t) \frac{\partial \omega}{\partial x_j} \right] + 2(1 - F_1) \frac{\rho \sigma_\omega k}{\omega} \frac{\partial k}{\partial x_i} \frac{\partial \omega}{\partial x_i}$$

Derivation of SST $k - \omega$ length scale limiter

SST $k - \omega$ [8]: RANS, eddy viscosity model, $\nu_t \approx \frac{k}{\omega}$, 2 eqn.— k, ω

- ω eqn.:
$$\frac{\partial(\rho\omega)}{\partial t} + \frac{\partial(\rho u_j \omega)}{\partial x_j} = \frac{\gamma}{\nu_t} \mathbf{P} - \beta \rho \omega^2 + \frac{\partial}{\partial x_j} \left[(\mu + \sigma_\omega \mu_t) \frac{\partial \omega}{\partial x_j} \right] + 2(1 - F_1) \frac{\rho \sigma_\omega \omega^2}{\omega} \frac{\partial k}{\partial x_i} \frac{\partial \omega}{\partial x_i}$$

Start from $k - \varepsilon$ eqns. w/ $C_{\varepsilon 1}^*$

Transform into $k - \omega$ eqns. w/ $C_{\varepsilon 1}^*$; find

- $\gamma \rightarrow \gamma^*$

Derivation of SST $k - \omega$ length scale limiter

SST $k - \omega$ [8]: RANS, eddy viscosity model, $\nu_t \approx \frac{k}{\omega}$, 2 eqn.— k, ω

- ω eqn.: $\frac{\partial(\rho\omega)}{\partial t} + \frac{\partial(\rho u_j \omega)}{\partial x_j} = \frac{\gamma}{\nu_t} \mathbf{P} - \beta \rho \omega^2 + \frac{\partial}{\partial x_j} \left[(\mu + \sigma_\omega \mu_t) \frac{\partial \omega}{\partial x_j} \right] + 2(1 - F_1) \frac{\rho \sigma_\omega \omega^2}{\omega} \frac{\partial k}{\partial x_i} \frac{\partial \omega}{\partial x_i}$

Start from $k - \varepsilon$ eqns. w/ $C_{\varepsilon 1}^*$

Transform into $k - \omega$ eqns. w/ $C_{\varepsilon 1}^*$; find

- $\gamma \rightarrow \gamma^*$

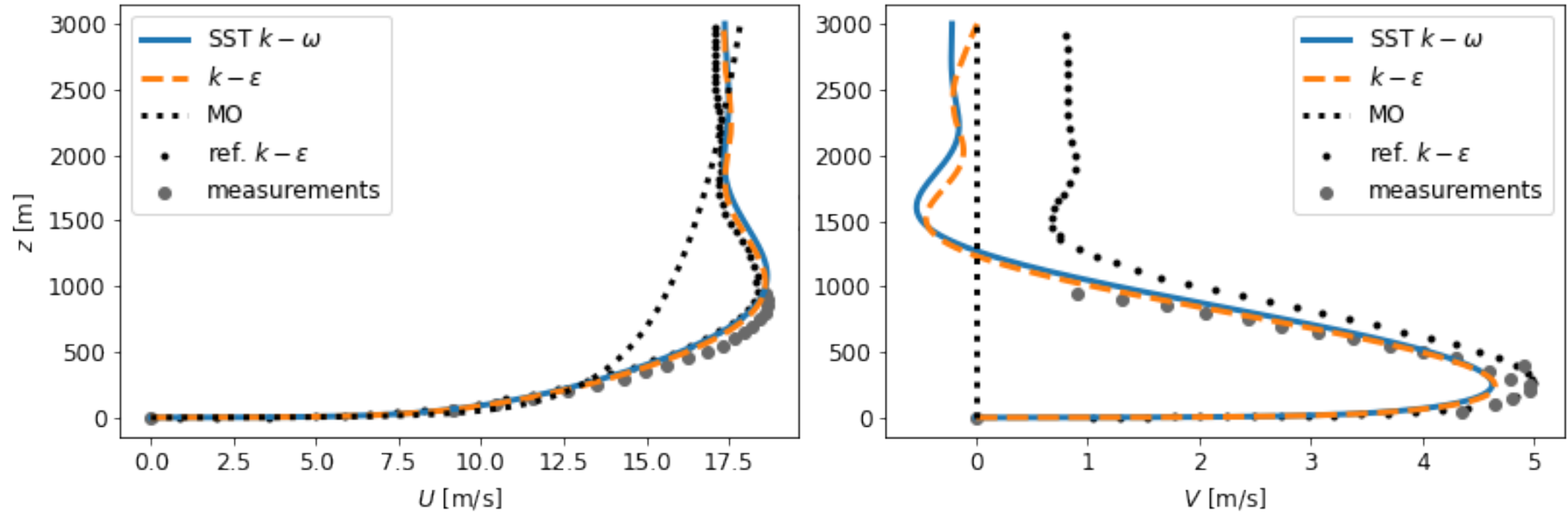
Apply to $k - \varepsilon$ and $k - \omega$ contributions to SST $k - \omega$:

- $\gamma_1^* = C_{\varepsilon 1}^* - 1$

- $\gamma_2^* = C_{\varepsilon 2}^* - 1$

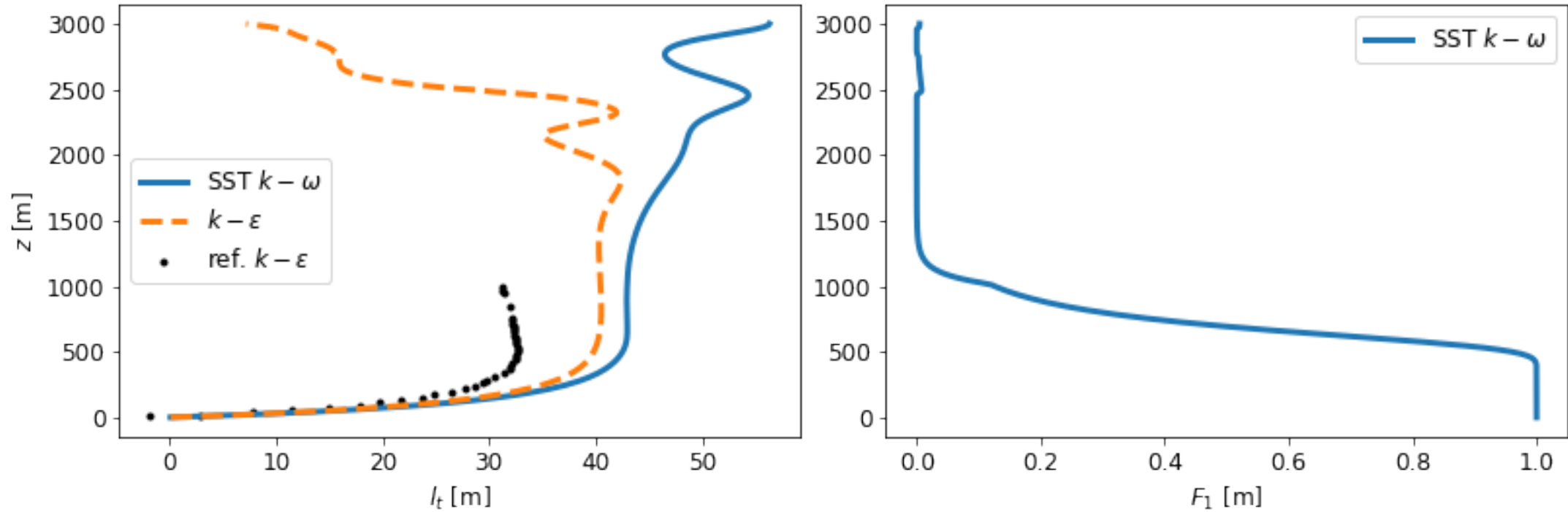
Blend: $\gamma^* = F \gamma_1^* + (1 - F) \gamma_2^*$

SST $k - \omega$ and $k - \varepsilon$ w/ limiter capture Coriolis effect



Reference $k - \varepsilon$ w/ limiter from [5] (also w/ wall function, extra diffusion)

Limiter stops growth of mixing length;
restores desired SST $k - \omega$ blending



Conclusions

- Showed baseline SST $k - \omega$ and $k - \varepsilon$ match MO similarity theory but don't fully capture Coriolis effect
- Derived mixing length limiter for SST $k - \omega$ using ideas from [5]
- Showed SST $k - \omega$ and $k - \varepsilon$ w/ limiter match measured ABL profile
- Showed $k - \varepsilon$ w/ limiter models this ABL case well w/out
 - Wall function, e.g. that in [5]
 - Diffusion term analogous to that in SST $k - \omega$, e.g. as in [5]

Future Work

- Hybrid RANS-LES (HRL) of ABL [9], specifically Active Model Split (AMS) [10], w/ SST $k - \omega$ w/ limiter
- SST $k - \omega$ and AMS of ABL including buoyancy
- AMS blade-resolved simulations of O(10) wind turbines in an ABL [1]

Works Cited

- [1] Sprague, M.A., Ananthan, S., Vijayakumar, G., Robinson, M., “ExaWind: A multi fidelity modeling and simulation environment for wind energy”, NAWEA/WindTech, 2019.
- [2] Menter, F.R., “Improved two-equation $k - \omega$ turbulence models for aerodynamic flows”, NASA Technical Memorandum 103975, 1992.
- [3] Wilcox, D.C., “Turbulence modeling for CFD”, DCW Industries, 2006.
- [4] Bautista, M.C., Dufresne, L., and Masson, C., “Hybrid Turbulence Models for Atmospheric Flow,” E3S Web of Conferences, 2015.
- [5] Koblitz, T., “CFD modeling of non-neutral atmospheric boundary layer conditions,” Ph.D. thesis, DTU, 2013.

Works Cited Cont.

- [6] Lettau, H., “A re-examination of the ‘Leipzig wind profile’ considering some relations between wind and turbulence in the frictional layer,” *Tellus*, 1950.
- [7] Launder, B.E, Spalding, D.B., “The numerical computation of turbulent flows”, *Computer Methods in Applied Mechanics and Engineering*, 1974.
- [8] Menter, F., Kuntz, M., Langtry, R., “Ten years of industrial experience with the SST turbulence model”, *Turbulence, Heat and Mass Transfer*, 2003.
- [9] Adcock, C., Henry de Frahan, M., Melvin, J., Vijayakumar, G., Ananthan, S., Iaccarino, G., Moser, R., Sprague, M., “Hybrid RANS-LES of the Atmospheric Boundary Layer for Wind Farm Simulations”, *AIAA SciTech*, 2021 (accepted.)
- [10] Haering, S. W., Oliver, T. A., and Moser, R. D., “Active Model Split Hybrid RANS/LES,” *arXiv*, 2020 (preprint.)

Acknowledgements

Christiane Adcock is supported in part by a graduate fellowship award from **Knight-Hennessy Scholars** at Stanford University.

This research was supported in part by the U.S. Department of Energy **Computational Science Graduate Fellowship** under grant DE SC0019323.

This research was supported by the **Exascale Computing Project** (17-SC-20-SC), a collaborative effort of the U.S. Department of Energy Office of Science and the National Nuclear Security Administration.

A portion of the research was performed using **computational resources** sponsored by the Department of Energy's Office of Energy Efficiency and Renewable Energy and located at the National Renewable Energy Laboratory.

This work was authored in part by the **National Renewable Energy Laboratory**, operated by Alliance for Sustainable Energy, LLC, for the U.S. Department of Energy (DOE) under Contract No. DE-AC36-08GO28308. Funding provided by U.S. Department of Energy Office of Science and National Nuclear Security Administration. The views expressed in the article do not necessarily represent the views of the DOE or the U.S. Government. The U.S. Government retains and the publisher, by accepting the article for publication, acknowledges that the U.S. Government retains a nonexclusive, paid-up, irrevocable, worldwide license to publish or reproduce the published form of this work, or allow others to do so, for U.S. Government purposes.

Questions

